UNIVERSITETET I STAVANGER

INSTITUTT FOR MATEMATIKK OG NATURVITENSKAP

FYS 610 Many-particle quantum mechanics

Suggested solutions, exercises for 7 April 2017

PROBLEM 22:

a) From the fundamental anticommutation relation, $\{\gamma^{\mu}, \gamma^{\nu}\} = 2g^{\mu\nu}\mathbb{1}_4$ it follows:

$$\left(\gamma^{0}\right)^{2}=\mathbb{1}_{4}; \qquad \left(\gamma^{i}\right)^{2}=-\mathbb{1}_{4}; \qquad \gamma^{\mu}\gamma^{j}=-\gamma^{j}\gamma^{\mu} \quad (\mu\neq j).$$

Thus from the definition of γ^5 :

$$\begin{split} \left(\gamma^{5}\right)^{2} &= -\gamma^{0} \gamma^{1} \gamma^{2} \gamma^{3} \gamma^{0} \gamma^{1} \gamma^{2} \gamma^{3} = -(-1)^{3} \gamma^{1} \gamma^{2} \gamma^{3} (\gamma^{0})^{2} \gamma^{1} \gamma^{2} \gamma^{3} = (-1)^{2} \gamma^{2} \gamma^{3} (\gamma^{1})^{2} \gamma^{2} \gamma^{3} \\ &= (-1)^{2} \gamma^{3} (\gamma^{2})^{2} \gamma^{3} = \mathbb{1}_{4} \,. \end{split}$$

b) From $\{\gamma^{\mu}, \gamma^{\nu}\} = 2g^{\mu\nu}$ it follows that $\gamma_{\mu}\gamma^{\mu} = g_{\mu}^{\ \mu} = \delta_{\mu\mu} = 4$, so:

$$p\gamma^{\mu} = p_{\nu}\gamma^{\nu}\gamma^{\mu} = p_{\nu}(2g^{\nu\mu} - \gamma^{\mu}\gamma^{\nu}) = 2p^{\mu} - \gamma^{\mu}p \iff \{p, \gamma^{\mu}\} = 2p^{\mu};$$
$$\gamma_{\mu}p\gamma^{\mu} = \gamma_{\mu}(2p^{\mu} - \gamma^{\mu}p) = 2p - 4p = -2p.$$

c) Using the previous result and $p^2 = p^2$, we have:

- d) From $\{\gamma^{\mu}, \gamma^{\nu}\} = 2g^{\mu\nu}$ it follows that γ^{μ} commutes with three of the gamma-matrix factors in the definition of γ^{5} and commutes with the fourth, γ^{μ} , so it anticommutes with γ^{5} .
- e) From the cyclical properties of the trace we have:

$$\operatorname{Tr}[\gamma^{\mu}\gamma^{\nu}] = \frac{1}{2}\operatorname{Tr}[\gamma^{\mu}\gamma^{\nu} + \gamma^{\nu}\gamma^{\mu}] = g^{\mu\nu}\operatorname{Tr}\mathbb{1}_{4} = 4g^{\mu\nu}.$$

Using this and the fundamental commutation relation we find:

$$\begin{aligned} \operatorname{Tr}[\gamma^{\alpha}\gamma^{\mu}\gamma^{\beta}\gamma^{\nu}] &= -\operatorname{Tr}[\gamma^{\alpha}\gamma^{\mu}\gamma^{\nu}\gamma^{\beta}] + 2g^{\beta\nu}\operatorname{Tr}[\gamma^{\mu}\gamma^{\alpha}] \\ &= -\operatorname{Tr}[\gamma^{\alpha}\gamma^{\mu}\gamma^{\nu}\gamma^{\beta}] + 8g^{\alpha\mu}g^{\beta\nu} \ . \end{aligned}$$

By repeated applications of this result we then have:

$$\begin{aligned} \operatorname{Tr}[\gamma^{\alpha}\gamma^{\mu}\gamma^{\beta}\gamma^{\nu}] &= -\operatorname{Tr}[\gamma^{\alpha}\gamma^{\mu}\gamma^{\nu}\gamma^{\beta}] + 8g^{\alpha\mu}g^{\beta\nu} \,. \\ &= \operatorname{Tr}[\gamma^{\alpha}\gamma^{\nu}\gamma^{\mu}\gamma^{\beta}] + 8(g^{\alpha\mu}g^{\beta\nu} - g^{\alpha\beta}g^{\mu\nu}) \\ &= -\operatorname{Tr}[\gamma^{\nu}\gamma^{\alpha}\gamma^{\mu}\gamma^{\beta}] + (g^{\alpha\mu}g^{\beta\nu} - g^{\alpha\beta}g^{\mu\nu} + g^{\alpha\nu}g^{\beta\mu}) \,. \end{aligned}$$

But the matrices in the trace in the last expression are just a cyclical rearrangement of those on the left hand side, so by the invariance of the trace under cyclical permutations, we find:

$$\mathrm{Tr}[\gamma^{\alpha}\gamma^{\mu}\gamma^{\beta}\gamma^{\nu}] = 4(g^{\alpha\mu}g^{\beta\nu} - g^{\alpha\beta}g^{\mu\nu} + g^{\alpha\nu}g^{\beta\mu}).$$

PROBLEM 23: Using the hint, $(\gamma^5)^2 = \mathbb{1}_4$, and the properties of the trace, we find

$$\operatorname{Tr}[\gamma^{\mu_1} \dots \gamma^{\mu_n}] = \operatorname{Tr}[\gamma^5 \gamma^{\mu_1} \dots \gamma^{\mu_n} \gamma^5] = (-1)^n \operatorname{Tr}[\gamma^{\mu_1} \dots \gamma^{\mu_n} (\gamma^5)^2]$$

SO

$$Tr[\gamma^{\mu_1} \dots \gamma^{\mu_n}] = 0,$$

if n is odd.

PROBLEM 24: We have:

$$AB = A_{\mu}B_{\nu}\frac{1}{2}\left(\{\gamma^{\mu},\gamma^{\nu}\} + [\gamma^{\mu},\gamma^{\nu}]\right) = A_{\mu}B_{\nu}\frac{1}{2}(2g^{\mu\nu} - 2i\sigma^{\mu\nu}) = A \cdot B \mathbb{1}_4 - iA_{\mu}B_{\nu}\sigma^{\mu\nu}.$$

Hence, if $[B_{\nu}, \gamma^{\mu}] = 0$, since $\operatorname{Tr} \sigma^{\mu\nu} = -2i \operatorname{Tr} [\gamma^{\mu} \gamma^{\nu} - \gamma^{\nu} \gamma^{\mu}] = 0$:

$$Tr[AB] = 4A \cdot B$$
.

PROBLEM 25: We note the following relations:

$$\gamma^{0\dagger} = \gamma^0, \qquad \gamma^{i\dagger} = -\gamma^i$$

which, together with the fundamental anticommutation relation yields

$$\gamma^{\mu\dagger} = \gamma^0 \gamma^\mu \gamma^0 \,.$$

Hence if u(p) is a Dirac spinor, satisfying $(\not p - m\mathbb{1}_4)u(p) = 0$ and $\bar{u}'(p) = u'^{\dagger}(p)\gamma_0$ satisfies the conjugated Dirac equation:

$$\bar{u}'(p)(\not p - m\mathbb{1}_4) = u'^{\dagger}(p)\gamma^0(p_{\mu}\gamma^{\mu} - m\mathbb{1}_4)\gamma^{0^2} = u'^{\dagger}(p)(p_{\mu}\gamma^{\mu\dagger}\gamma_0 - m\mathbb{1}_4)\gamma_0$$
$$= \left[\gamma_0(p_{\mu}\gamma^{\mu}\gamma_0 - m\mathbb{1}_4)u'(p)\right]^{\dagger} = 0$$

If a^{μ} is an arbitrary 4-vector, we then have, using the result in the previous problem:

$$\begin{split} 0 &= \bar{u}'(q)(\not q - m\mathbb{1}_4) \not a u(p) + \bar{u}(q) \not a (\not p - m\mathbb{1}_4) u(p) \\ &= \bar{u}'(q) \left[-2m \, \gamma^\mu a_\mu + (p^\mu + q^\mu) a_\mu + \mathrm{i} (q_\nu - p_\nu) \sigma^{\mu\nu} a_\mu \right] u(p) \end{split}$$

But this equation is linear in a^{μ} , and if it is satisfied for all values of a^{μ} , the coefficient of a_{μ} must vanish. Equivalently, we may differentiate the equation with respect to a_{μ} . Hence we have proven the Gordon decomposition:

$$\bar{u'}(q)\gamma^{\mu}u(p) = \bar{u}'(q)\left[\frac{p^{\mu}+q^{\mu}}{2m} + \mathrm{i}\frac{q^{\nu}-p^{\nu}}{2m}\sigma^{\mu\nu}\right]u(p)\,.$$

Problem 26 [Not yet available]

PROBLEM 27 [Not yet available]